

# Thermodynamic optimization of the Yb–Sn system

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**Abstract** The thermodynamic optimization of the Yb–Sn binary system was carried out with the help of the CALPHAD (CALculation of PHAse Diagram) method.  $\text{Yb}_2\text{Sn}$ ,  $\alpha$ - $\beta$  $\text{Yb}_5\text{Sn}_3$ ,  $\text{Yb}_5\text{Sn}_4$ ,  $\text{YbSn}$ ,  $\text{Yb}_3\text{Sn}_5$ , and  $\text{YbSn}_3$  have been treated as stoichiometric compounds, while a sublattice model has been used for the description of the liquid, BCC, BCT, and FCC phases. The calculations based on the thermodynamic modeling are in good agreement with the phase diagram data and experimental thermodynamic values.

**Keywords** Yb–Sn system · Phase diagram · Thermodynamic optimization · Calphad method

## Introduction

This study is a part of a thermodynamic investigation of the R–Sn systems (R=La) [1] which is intended to give a better understanding of the constitutional properties and potential technological applications of these alloys.

This study deals with an assessment of the thermodynamic description of the Yb–Sn system using the

CALPHAD technique [2]. The thermodynamic parameters involved in the models are optimized from the experimental thermodynamic and phase diagram data.

## Review of experimental data

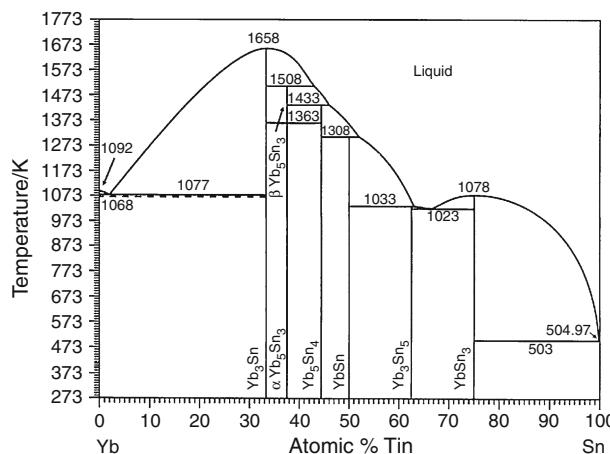
The assessed Yb–Sn phase diagram by Palenzona and Cirafici [3] is based on the work of Palenzona and Cirafici [4]. According to this assessment, five intermetallic compounds were reported ( $\text{YbSn}_3$ ,  $\text{YbSn}$ ,  $\text{Yb}_5\text{Sn}_4$ ,  $\text{Yb}_5\text{Sn}_3$ , and  $\text{Yb}_2\text{Sn}$ ). For the  $\text{YbSn}_3$  and  $\text{Yb}_2\text{Sn}$  compounds, a congruent melting was determined, respectively, at 1078 and 1658 K, while for the  $\text{YbSn}$ ,  $\text{Yb}_5\text{Sn}_4$ , and  $\text{Yb}_5\text{Sn}_4$  ( $\alpha$  and  $\beta$ ) compounds, a peritectic formation was reported, respectively, at 1308, 1433, and 1508 K. A polymorphic transition ( $\alpha \leftrightarrow \beta$ ) was determined at 1363 K for  $\text{Yb}_5\text{Sn}_4$ . In addition, two more phases were reported in recent papers, namely  $\text{Yb}_3\text{Sn}_5$  [5] (see Fig. 1) and  $\text{Yb}_{36}\text{Sn}_{23}$  [6]. The investigation of the Yb–Sn phase diagram is severely hampered by the high oxidability of the alloys in the central part of the system [3]. Moreover, a careful investigation on two of the phases reported in the assessment, namely  $\text{Yb}_3\text{Sn}_5$  and  $\text{Yb}_5\text{Sn}_4$ , has shown that in the Yb–Sn system, these phases are stabilized by hydrogen impurity [7]. The crystal structures of the various phases of the Yb–Sn system are reported in Table 1.

Many authors [8–10] determined the enthalpies of formation for the intermetallic compounds of the Yb–Sn system; Palenzona [8] measured the standard enthalpy of formation of  $\text{YbSn}_3$  using the dynamic differential calorimetry method. Colinet et al. [9] calculated by a semi-empirical model the enthalpies of formation of  $\text{YbSn}_3$  [10] determined the enthalpies of formation for two intermetallic compounds  $\text{Yb}_5\text{Sn}_4$  and  $\text{Yb}_5\text{Sn}_3$ .

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**Fig. 1** The Yb–Sn system. Phase diagram revised in the composition range between 55 and 70 at.%Sn by Manfrinetti et al. [5]

## Thermodynamic models

### Pure elements

The Gibbs energy function

$$G_i^\phi(T) = {}^0G_i^\phi - H_i^{\text{SER}}(298.15) \quad (1)$$

(298.15 K) for the element  $i$  ( $i = \text{Yb, Sn}$ ) in the phase  $\Phi$  ( $\Phi = \text{Liquid, BCC\_A2, HCP\_A3, BCT\_A5, or FCC\_A1}$ ) is described by an equation of the following form:

$$G_i^\phi(T) = a + bT + cT \ln T + dT^{-2} + eT^3 + fT^7 + gT^{-9}, \quad (2)$$

where  $H_i^{\text{SER}}$  (298.15 K) is the molar enthalpy of the element  $i$  at 298.15 K in its standard element reference (SER) state, FCC\_A1 for Yb and BCT\_A5 for Sn.

In this article, the Gibbs energy functions are taken from the SGTE compilation of Dinsdale [11].

### Solution phases

The solution phases [ $(\gamma\text{Yb}), \beta(\text{Yb}), \alpha(\text{Yb}), \beta(\text{Sn}), \text{ and liquid}$ ] were modeled as substitutional solutions according to the polynomial Redlich–Kister model [12]. The Gibbs energy of 1 mol of formula unit of phase  $\phi$  is expressed as the sum of the reference part  ${}^{\text{ref}}G^\phi$ , the ideal part  ${}^{\text{id}}G^\phi$ , and the excess part  ${}^{\text{xs}}G^\phi$ :

$$G_m^\phi = {}^{\text{ref}}G^\phi + {}^{\text{id}}G^\phi + {}^{\text{xs}}G^\phi. \quad (3)$$

As used in the Thermo-Calc software [13]:

$$\begin{aligned} {}^{\text{ref}}G^\phi(T) &= \left({}^0G_{\text{Yb}}^\phi(T) - H_{\text{Yb}}^{\text{SER}}(298.15)\right)x_{\text{Yb}} \\ &+ \left({}^0G_{\text{Sn}}^\phi(T) - H_{\text{Sn}}^{\text{SER}}(298.15)\right)x_{\text{Sn}} \end{aligned} \quad (4)$$

$${}^{\text{id}}G^\phi = RT(x_{\text{Sn}} \ln x_{\text{Sn}} + x_{\text{Yb}} \ln x_{\text{Yb}}) \quad (5)$$

where  $R$  is the gas constant,  $T$  the temperature, in Kelvin,  $x_{\text{Sn}}$  and  $x_{\text{Yb}}$  are the mole fraction of elements Sn and Yb, respectively.

The excess terms of all the phases were modeled by the Redlich–Kister [12] formula.

$$\begin{aligned} {}^{\text{xs}}G_m^\phi(T) &= x_{\text{Sn}}x_{\text{Yb}} \left( {}^0L_{\text{Sn,Yb}}^\phi(T) + {}^1L_{\text{Sn,Yb}}^\phi(T)(x_{\text{Sn}} - x_{\text{Yb}}) \right. \\ &\quad \left. + {}^2L_{\text{Sn,Yb}}^\phi(T)(x_{\text{Sn}} - x_{\text{Yb}})^2 + \dots \right) \end{aligned} \quad (6)$$

with

$${}^iL_{\text{Sn,Yb}}^\phi(T) = a_i + b_iT \quad (7)$$

**Table 1** Symbols and crystal structures of the stable solid phases in the (Yb–Sn) alloys from [3, 5]

Diagram symbol	Composition/ at.% Yb	Symbol used in Thermo-Calc data file	Pearson symbol	Space group	Struktur- bericht designation	Prototype
$\beta\text{Sn}$	0	BCT	$tI4$	$I4_1/\text{amd}$	$A5$	$\beta\text{Sn}$
$\alpha\text{Sn}$	0	Diamond	$cF8$	$Fd\bar{3}m$	$A5$	C (diamond)
$\text{YbSn}_3$	25	$\text{YbSn}_3$	$cP4$	$Fm\bar{3}m$	$L1_2$	$\text{AuCu}_3$
$\text{YbSn}$	50	$\text{YbSn}$	$tP2$	$P4/mmm$	$L1_0$	$\text{AuCu-I}$
$\text{Yb}_5\text{Sn}_4$	55.6	$\text{Yb}_5\text{Sn}_4$	$oP36$	$Pnma$	—	$\text{Gd}_5\text{Si}_4$
$\beta\text{Yb}_5\text{Sn}_3$	62.5	$\text{Yb}_5\text{Sn}_3\text{-HT}$	$tl32$	$I4/mcm$	$D8_l$	$\text{Cr}_5\text{B}_3$
$\alpha\text{Yb}_5\text{Sn}_3$	62.5	$\alpha\text{Yb}_5\text{Sn}_3\text{-LT}$	$hP16$	$P6_3/mcm$	$D8_8$	$\text{Mn}_5\text{Si}_3$
$\text{Yb}_2\text{Sn}$	66.7	$\text{Yb}_2\text{Sn}$	$hP6$	$P6_3/mcm$	$B8_2$	$\text{Ni}_2\text{In}$
$\gamma\text{Yb}$	100	BCC	$cl2$	$Im\bar{3}m$	$A2$	W
$\beta\text{Yb}$	100	FCC	$cF4$	$Fm\bar{3}m$	$A1$	Cu
$\alpha\text{Yb}$	100	HCP	$hP2$	$P6_3/mmc$	$A3$	Mg

where  ${}^iL_{\text{Sn}, \text{Yb}}^\phi(T)$  is the  $i$ th interaction parameter between the elements Sn and Yb, which is evaluated in the presented work,  $a_i$  and  $b_i$  are the coefficients to be optimized.

### Stoichiometric compounds

The Gibbs energy of the stoichiometric compound  ${}^0G_{\text{ApBq}}$  is expressed as follows:

$${}^0G_{\text{ApBq}} = \frac{p}{p+q} {}^0G_A + \frac{q}{p+q} {}^0G_B + a + bT, \quad (8)$$

where  ${}^0G_A$  and  ${}^0G_B$  are the Gibbs energy of the pure elements Sn and Yb, respectively,  $a$  and  $b$  are parameters to be determined.

### Results and discussions

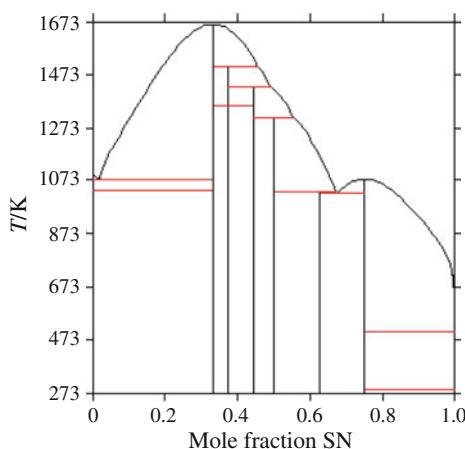
The evaluation of the thermodynamic parameters has been carried out using the PARROT [14] module in the Thermo-Calc code [13], takes various types of experimental data for the optimisation process. Data for the pure elements have been taken from the SGTE compilation [11]. Most of the experimental data are mentioned in “[Review of experimental data](#)” section. In the present parameter optimization procedure, we first imposed the conditions  $d^2G/dx^2 > 0$  for modeling the liquid phase using the phase boundary data reported by [5] and thermodynamic data. We then modeled the parameters of the intermetallic phases beginning with the congruent melting ( $\text{Liq} \rightleftharpoons \text{Yb}_2\text{Sn}$ ). Finally, all the parameters of the different phases were simultaneously optimized with the experimental data mentioned above by imposing additional constraints to avoid the appearance of unwanted inverted miscibility gap in the liquid phase during the phase diagram calculation as recommended in [15–17]. All the parameters were evaluated and listed in Table 2.

The calculated phase diagram is shown in Fig. 2. The experimental and calculated temperatures of the invariant reactions are compared in Table 3. They are in very good agreement. The measured and calculated enthalpies of formation of the intermetallic compounds are presented in Fig. 3. They are in reasonable agreement.

The optimized temperature of the polymorphic transition ( $1090\ ^\circ\text{C}$ ) of  $\alpha\text{Yb}_5\text{Sn}_3 \rightleftharpoons \beta\text{Yb}_5\text{Sn}_3$  chosen in [4] is in agreement with the calculated one ( $1087\ ^\circ\text{C}$ ). The enthalpies of formation of the two polytypes,  $\alpha\text{Yb}_5\text{Sn}_3$  and  $\alpha\text{Yb}_5\text{Sn}_3 \rightleftharpoons \beta\text{Yb}_5\text{Sn}_3$ , are very similar and confirm the conclusion by Stein et al. [18] concerning the small structural energies of the different Laves phases polytypes. As mentioned in [17], in order to check that the optimized thermodynamic parameters of the intermetallic compounds are satisfactory, we verified that when the

**Table 2** The optimised thermodynamic parameters of the Yb–Sn system

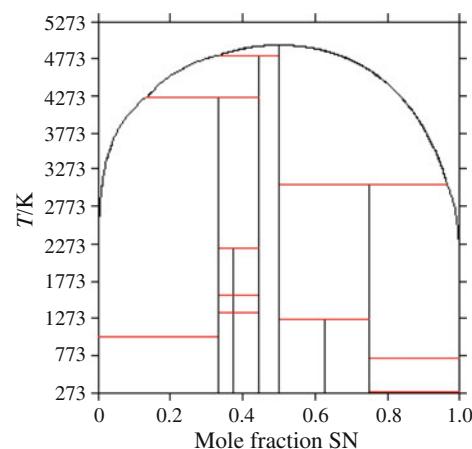
Phase	Liquid	Thermodynamic models	Parameters in SI
BCC_A2	(Yb, Sn) <sub>1</sub> (Va) <sub>3</sub>	${}^0L_{\text{Yb,Sn}}^{\text{Liq}} = -155596.489 - 1.401083T$	No excess term
FCC_A1	(Yb, Sn) <sub>1</sub> (Va) <sub>3</sub>	${}^1L_{\text{Yb,Sn}}^{\text{Liq}} = -3736.52101$	$G_{\text{Yb,Sn}}^{\text{YbSn}_3} - 0.25 {}^{298}H_{\text{Yb}}^{\text{FCC\_A1}} - 0.75 {}^{298}H_{\text{Sn}}^{\text{BCT\_A5}} = 0.25 {}^{298}G_{\text{Yb}}^{\text{FCC\_A1}} + 0.75 {}^{298}G_{\text{Sn}}^{\text{BCT\_A5}} - 41797.2855 + 2.55025052T$
YbSn3	(Yb) <sub>0.25</sub> :(Sn) <sub>0.75</sub>	${}^0L_{\text{Yb,Sn}}^{\text{Liq}} = -3736.52101$	$G_{\text{Yb,Sn}}^{\text{YbSn}_3} - 0.5 {}^{298}H_{\text{Yb}}^{\text{FCC\_A1}} - 0.5 {}^{298}H_{\text{Sn}}^{\text{BCT\_A5}} = 0.5 {}^{298}G_{\text{Yb}}^{\text{FCC\_A1}} + 0.5 {}^{298}G_{\text{Sn}}^{\text{BCT\_A5}} - 55077.4602 + 0.646652440T$
YbSn	(Yb) <sub>0.5</sub> :(Sn) <sub>0.5</sub>	${}^2L_{\text{Yb,Sn}}^{\text{Liq}} = 36185.9184$	$G_{\text{Yb,Sn}}^{\text{YbSn}_4} - 0.556 {}^{298}H_{\text{Yb}}^{\text{FCC\_A1}} - 0.444 {}^{298}H_{\text{Sn}}^{\text{BCT\_A5}} = 0.556 {}^{298}G_{\text{Yb}}^{\text{FCC\_A1}} + 0.444 {}^{298}G_{\text{Sn}}^{\text{BCT\_A5}} - 56787.7363 + 1.37672739T$
Yb <sub>5</sub> Sn <sub>4</sub>	(Yb) <sub>0.556</sub> :(Sn) <sub>0.444</sub>		$G_{\text{Yb,Sn}}^{\text{YbSn}_3} - 0.625 {}^{298}H_{\text{Yb}}^{\text{FCC\_A1}} - 0.375 {}^{298}H_{\text{Sn}}^{\text{BCT\_A5}} = 0.625 {}^{298}G_{\text{Yb}}^{\text{FCC\_A1}} + 0.375 {}^{298}G_{\text{Sn}}^{\text{BCT\_A5}} - 57602.0929 + 0.337700156T$
$\beta\text{Yb}_5\text{Sn}_3$	(Yb) <sub>0.625</sub> :(Sn) <sub>0.375</sub>		$G_{\text{Yb,Sn}}^{\text{YbSn}_3} - 0.625 {}^{298}H_{\text{Yb}}^{\text{FCC\_A1}} - 0.375 {}^{298}H_{\text{Sn}}^{\text{BCT\_A5}} = 0.625 {}^{298}G_{\text{Yb}}^{\text{FCC\_A1}} + 0.375 {}^{298}G_{\text{Sn}}^{\text{BCT\_A5}} - 58406.1714 + 3.21533668T$
$\alpha\text{Yb}_5\text{Sn}_3$	(Yb) <sub>0.625</sub> :(Sn) <sub>0.375</sub>		$G_{\text{Yb,Sn}}^{\text{YbSn}_3} - 0.667 {}^{298}H_{\text{Yb}}^{\text{FCC\_A1}} - 0.333 {}^{298}H_{\text{Sn}}^{\text{BCT\_A5}} = 0.667 {}^{298}G_{\text{Yb}}^{\text{FCC\_A1}} + 0.333 {}^{298}G_{\text{Sn}}^{\text{BCT\_A5}} - 57560.2159 + 3.50846694T$
Yb <sub>2</sub> Sn	(Yb) <sub>0.667</sub> :(Sn) <sub>0.333</sub>		(Va) for vacancy



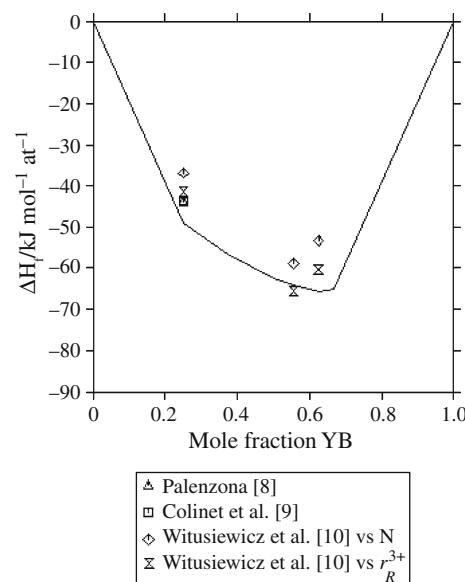
**Fig. 2** Calculated Yb–Sn phase diagram

**Table 3** Invariant reactions in the Yb–Sn system

Reaction	Refs. [3, 5]		This work	
	T/K	$x_{\text{Liq}}$ /at.% Yb	T/K	$x_{\text{Liq}}$ /at.% Yb
$\text{Liq} \rightleftharpoons \beta\text{Sn} + \text{YbSn}_3$	503	0.005	505	0.000045
$\text{Liq} \rightleftharpoons \text{YbSn}_3$	1078		1076	
$\text{Liq} \rightleftharpoons \text{YbSn}_3 + \text{Yb}_3\text{Sn}_5$	1023	0.33	1026	0.324
$\text{Liq} \rightleftharpoons \text{Yb}_3\text{Sn}_5 + \text{YbSn}$	1033	0.375	1032	0.33
$\text{Liq} + \text{Yb}_5\text{Sn}_4 \rightleftharpoons \text{YbSn}$	1308	0.48	1311	0.44
$\text{Liq} + \beta\text{Yb}_5\text{Sn}_3 \rightleftharpoons \text{Yb}_5\text{Sn}_4$	1433	0.53	1428	0.507
$\text{Liq} + \text{Yb}_2\text{Sn} \rightleftharpoons \beta\text{Yb}_5\text{Sn}_3$	1508	0.57	1505	0.543
$\alpha\text{Yb}_5\text{Sn}_3 \rightleftharpoons \beta\text{Yb}_5\text{Sn}_3$	1363	–	1360	
$\text{Liq} \rightleftharpoons \text{Yb}_2\text{Sn}$	1658	0.667	1663	0.667
$\text{Liq} \rightleftharpoons \text{Yb}_2\text{Sn} + \gamma\text{Yb}$	1077	0.98	1074	0.98



**Fig. 4** Calculated La–Pb phase diagram when the liquid phase is suspended



**Fig. 3** Calculated and measured enthalpies of formation of the intermetallic compounds

liquid phase is suspended during the calculation of the Yb–Sn phase diagram, the stoichiometric phases disappear at high temperatures, the terminal solid solutions and a two-phase domain existing between them are calculated, Fig. 4. In the present case, the entropies of the intermetallic compounds are either negative ( $s = -b$  optimized positive value; Table 2) and lower than 3.5 J/mol K. It will be noted that only the (Yb–Sn) BCC\_A2 solid solution is calculated on the whole Yb composition range and not the BCT\_A5 ( $\beta\text{Sn}$ ) one in the Sn-rich region. This is due to the power series in terms of temperature for the Sn element in the BCC\_A2 state which becomes metastable compared with the BCT\_A5 state at higher temperatures [11].

## Conclusions

A consistent set of thermodynamic parameters of the different phases of the Yb–Sn binary system has been optimized.

The computed values are in good agreement with the experimental data. We verified that no unwanted inverted miscibility gap was calculated for the solution phases.

Further thermodynamic determinations in particular for the liquid phase will be necessary to improve the assessment.

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